

# Gap Probabilities in Two-Matrix Models

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GERNOT AKEMANN



work with:

+ P. DAMGAARD: arXiv:0803.1171 (JHEP) & in preparation

+ J. OSBORN + K. SPLITTORFF: hep-th/0609059v2 (Nucl.Phys.B)

# Outline

1. **Two Matrix Model** (2MM)
2. **Physics** Motivation: **QCD**
3. **Objects:**  
densities, gaps, individual eigenvalues
4. **Solution for densities**
5. Expansion of gaps & eigenvalues in densities
6. **Exact results for gaps & eigenvalues**
7. **Large- $N$  & Examples**
8. Remarks and open problems

## Chiral Two-Matrix Model

$$\mathcal{Z} = \int d\Phi d\Psi \prod_{f=1}^{N_1} \det[i\mathcal{D}_1 + m_f] \prod_{g=1}^{N_2} \det[i\mathcal{D}_2 + n_g] e^{-\text{Tr}(\Phi\Phi^\dagger + \Psi^\dagger\Psi)}$$

- we are interested in real non-neg. eigenvalues of

$$\mathcal{D}_{1,2} \equiv \begin{pmatrix} 0 & \Phi + \mu_{1,2}\Psi \\ \Phi^\dagger + \mu_{1,2}\Psi^\dagger & 0 \end{pmatrix}$$

$\Phi, \Psi$  are matrices of sizes  $N \times (N + \nu) \in \mathbb{C}$

- change of variables  $\Phi_j \equiv \Phi + \mu_j\Psi$ ,  $j = 1, 2$  couples them:

$$\mathcal{Z} = \int d\Phi_1 d\Phi_2 \prod_{f=1}^{N_1} \det[\mathcal{D}_1 + m_f] \prod_{g=1}^{N_2} \det[\mathcal{D}_2 + n_g] e^{-\text{Tr}(c_1\Phi_1\Phi_1^\dagger + c_2\Phi_2\Phi_2^\dagger - d(\Phi_1\Phi_2^\dagger + \Phi_1^\dagger\Phi_2))}$$

- Unitary dof  $\Phi_1 = U_1 X V_1$  can be eliminated using a group integral (HCIZ for non-chiral)
- $\Phi_i$  Hermitean  $N \times N$  for non-chiral [Ercolani,McLaughlin01]

## Physics Motivation

- low energy **QCD**  $\longrightarrow$  chiral Perturbation Theory ( $\chi$ PT)  
 $\longrightarrow$  const. Pion fields = Matrix Model ( $N \rightarrow \infty$ )

[Shuryak, Verbaarschot 93; ...]

$$\mathcal{Z}_{\epsilon\chi PT} = \int_{U(N_f)} dU_0 \det[U_0]^\nu e^{\frac{1}{4}VF^2\mu^2 \text{Tr}[\Gamma, U_0][\Gamma, U_0^\dagger] + \frac{1}{2}\Sigma V \text{Tr}(M(U_0 + U_0^\dagger))}$$

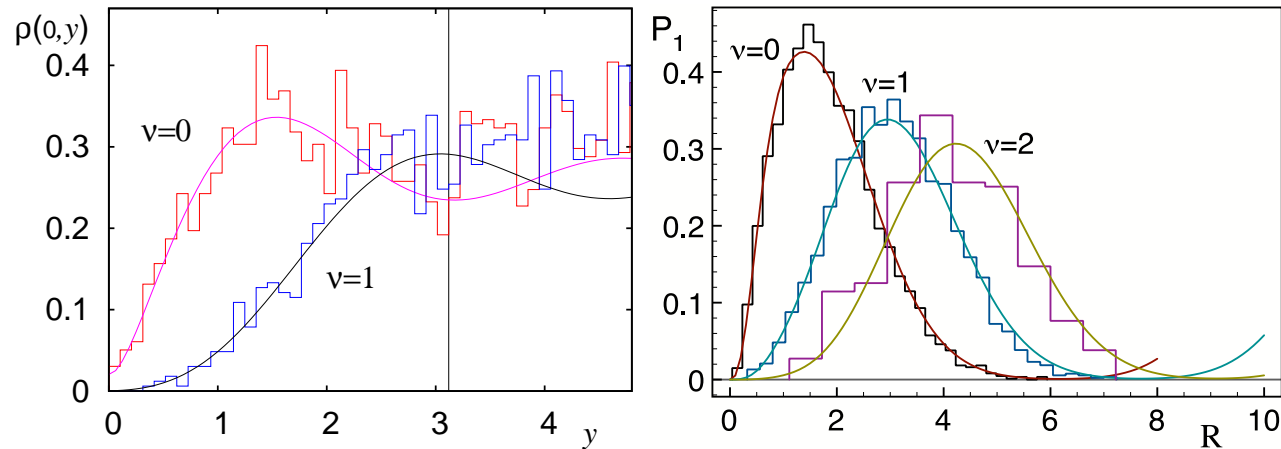
- effective  $\chi$ PT has 2 LEC's:  $F$  Pion decay constant &  $\Sigma$  chiral condensate
- determine by comparing *analytical* MM predictions to Lattice **QCD**  $\longrightarrow$  need source  $\mu$  to couple to  $F$
- easiest: imaginary isospin chemical potential  $i\mu$   
 $\longrightarrow$  keeps  $\mathcal{D}(\pm i\mu)$  spectra real!

[Damgaard, Heller, Splittorff, Svetitzky, Toublan 05/06]

- Individual eigenvalues: num. easiest [G.A., Damgaard 08]

## Illustration Density vs. '1st' Eigenvalue on $\mathbb{C}$

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- cut through density (left) vs. Lattice data [Bloch,Wettig'06]
- 1st eigenvalue (right) radially integrated **vs same data**  
[+ G.A., Shifrin 07]

# Eigenvalue Representation

- two sets of real eigenvalues:  $\{x\}$ ,  $\{y\}$

diagonalise  $\Phi_1 = U_1 X V_1$ ,  $\Phi_2 = U_2 Y V_2$  + use integral:

[Guhr, Wettig; Jackson, Sener, Verbaarschot 96]

$$\int dU dV \exp [N d \Re \text{Tr}(V X U Y)] = \frac{\det_{1 \leq i, j \leq N} [I_\nu(N dx_i y_j)]}{\prod_{i=1}^N (x_i y_i)^\nu \Delta_N(\{x^2\}) \Delta_N(\{y^2\})}$$

- chiral 2MM:

$$\mathcal{Z} = \prod_{i=1}^N \int_0^\infty dx_i dy_i (x_i y_i)^{\nu+1} \prod_{f=1}^{N_1} (x_i^2 + m_f^2) \prod_{g=1}^{N_2} (y_i^2 + n_g^2) \\ \times \Delta_N(x^2) \Delta_N(y^2) \det_{k,l} [I_\nu(2dN x_k y_l)] e^{-N \sum_i^N (c_1 x_i^2 + c_2 y_i^2)}$$

- non-Gaussian!! (non-chiral:  $I_\nu \rightarrow e$ )
- inserted  $\mathcal{D}$  operators = ave characteristic polynomials

## Objects of Study

- Define jpdf:  $\mathcal{Z} \equiv \prod_{i=1}^N \int_0^\infty dx_i dy_i \mathcal{P}_{jpdf}$

- **densities:**

$$R_{k,l}(x_1, \dots, x_k; y_1, \dots, y_l) \sim \prod_{i=k+1}^N \int_0^\infty dx_i \prod_{j=l+1}^N \int_0^\infty dy_j \mathcal{P}_{jpdf}$$

- **gap probabilities:**

$$E_{k,l}(s, t) \sim \int_0^s dx_1 \dots dx_k \int_s^\infty dx_{k+1} \dots dx_N \int_0^t dx_1 \dots dx_l \int_t^\infty dy_{l+1} \dots dy_N \mathcal{P}_{jpdf}$$

- **individual eigenvalues:**  $x_k = s$  &  $y_l = t$

$$p_{k,l}(s, t) \sim \int_0^s dx_1 \dots dx_{k-1} \int_s^\infty dx_{k+1} \dots dx_N \\ \times \int_0^t dy_1 \dots dy_{l-1} \int_t^\infty dy_{l+1} \dots dy_N \mathcal{P}_{jpdf}(\dots, x_k = s, \dots, y_l = t, \dots)$$

- not independent: e.g.  $\partial_s \partial_t E_{0,0}(s, t) = p_{1,1}(s, t)$  for smallest

## Solution for all Densities

- **bi-orthogonal polynomials**  $\delta_{nk} = \int_0^\infty dx dy w^{(N_1, N_2)}(x, y) P_n(x^2) Q_k(y^2)$

$$\text{weight } w^{(N_1, N_2)}(x, y) = (xy)^{\nu+1} \prod_{f=1}^{N_1} (x^2 + m_f^2) \prod_g^{N_2} (y^2 + m_g^2) I_\nu(2dNxy) e^{-N(c_1x^2 + c_2y^2)}$$

- in terms of **4 kernels**: using [Eynard, Mehta 98]

$$K_N = \sum_k^{N-1} P_k Q_k, \quad H_N = \sum_k^{N-1} \hat{\chi}_k P_k, \quad \hat{H}_N = \sum_k^{N-1} \chi_k Q_k, \quad M_N = \sum_k^{N-1} \hat{\chi}_k \chi_k$$

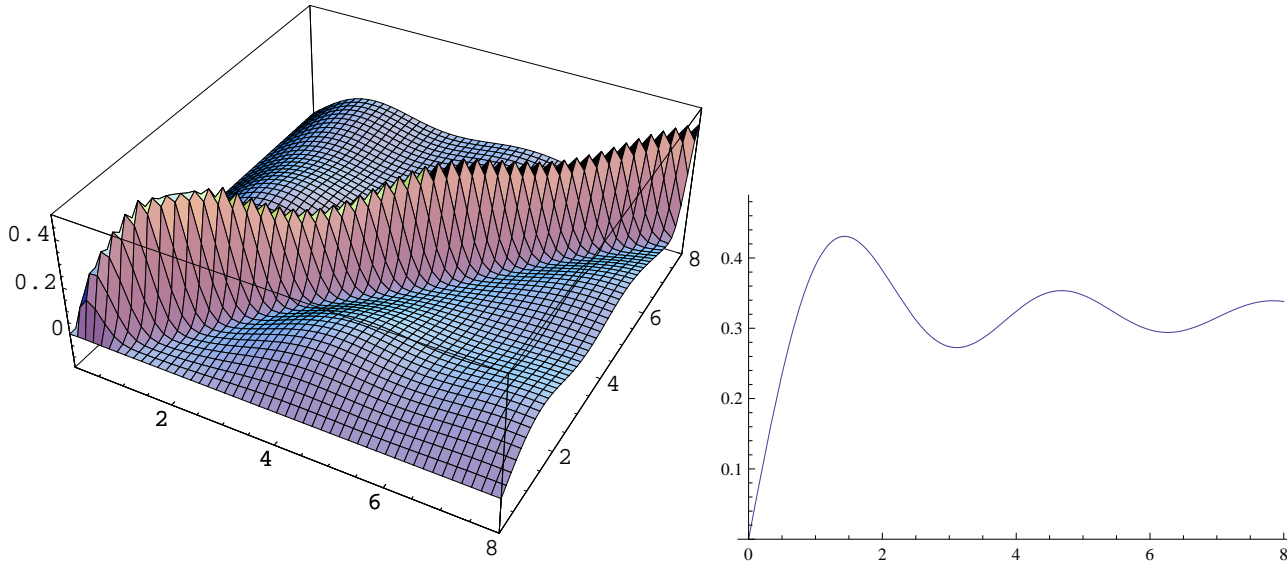
containing **integral trafo**  $\chi_k(x) = \int_0^\infty dy w(x, y) P_k(y)$  & ditto for  $Q_k$

⇒ all density correlation functions known explicitly

$$R_{n,k}(\{x\}_n; \{y\}_k) = \det_{\substack{1 \leq i_1, i_2 \leq n; \\ 1 \leq j_1, j_2 \leq k}} \begin{bmatrix} H_N(x_{i_1}, x_{i_2}) & M_N(x_{i_1}, y_{j_2}) - w^{(N_1, N_2)}(x_{i_1}, y_{j_2}) \\ K_N(y_{j_1}, x_{i_2}) & \hat{H}_N(y_{j_1}, y_{j_2}) \end{bmatrix}$$

## Example Density: 2MM vs. 1MM

example quenched ( $N_1 = N_2 = 0$ ):  $P_k, Q_k \sim L_k^\nu$  and  $\chi_k, \hat{\chi}_k \sim e^{-x^2} L_k^\nu$  **Laguerre**



- 1MM  $N \rightarrow \infty$ : **Bessel**  $\rho_1(x) = \frac{x}{2}(J_\nu(x)^2 - J_{\nu+1}(x)J_{\nu-1}(x))$
- 2MM  $N \rightarrow \infty$ : **generalised Bessel**  $\alpha^2 = 2N\delta\mu^2$

$$\rho_{1,1}(x, y) = \rho_1(x)\rho_1(y) - xy \mathcal{K}^+(x, y) \left( \mathcal{K}^-(x, y) - e^{-(x^2+y^2)/\alpha^2} I_\nu(xy/\alpha^2) \right)$$

$$\mathcal{K}^\pm(x, y) \equiv \int_0^1 dt e^{\pm t^2 \alpha^2} J_\nu(tx) J_\nu(ty)$$

## Densities – Gaps – Eigenvalues: Expansion

- all densities  $\longrightarrow$  all gap prob's  $\longrightarrow$  all indiv. ev.

$$E_{k,l}(s, t) = \sum_{i=0}^{N-k} \sum_{j=0}^{N-l} \frac{(-)^{i+j}}{i!j!} \int_0^s dx_1 \dots dx_{k+i} \int_0^t dy_1 \dots dy_{l+j} R_{k+i,l+j}(\{x\}_{k+i}, \{y\}_{l+j})$$

- example  $p_{1,1} = \partial\partial E_{0,0}$ :

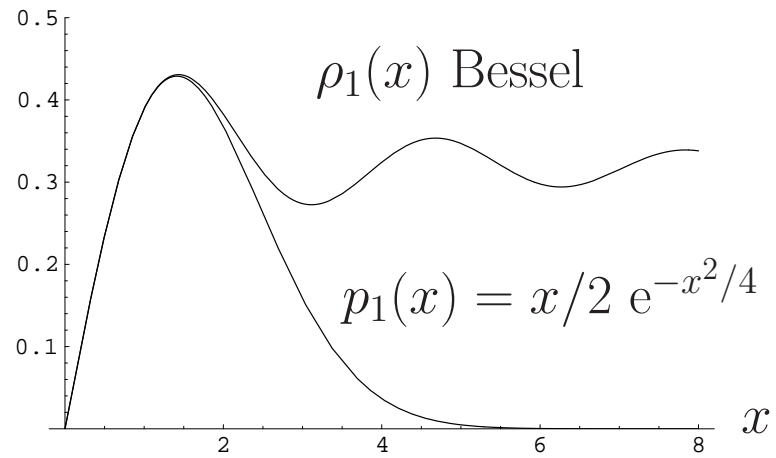
$$p_{1,1}(s, t) = R_{1,1}(s, t) - \int_0^s dx R_{2,1}(x, s, t) - \int_0^t dy R_{1,2}(s, t, y) + \dots$$

- same expansion for 1MM:

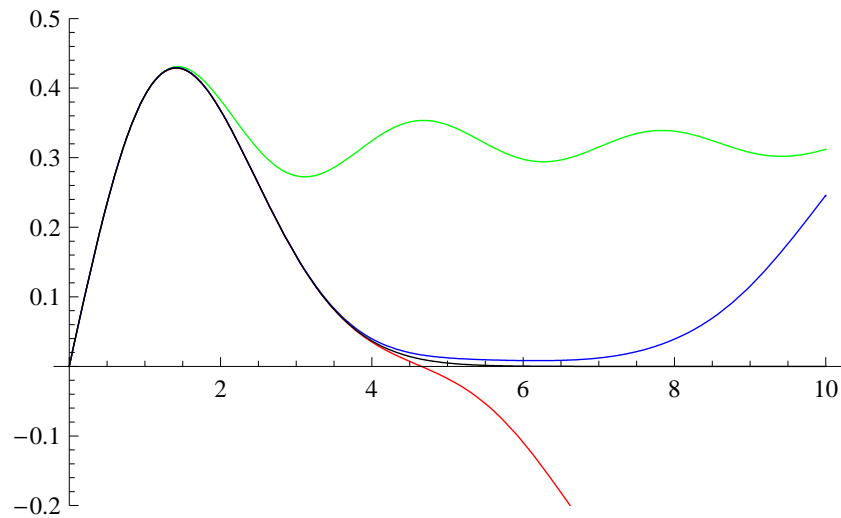
**1st eigenvalue = density –  $\int$  2-pt density +...**

# Expansion in 1MM

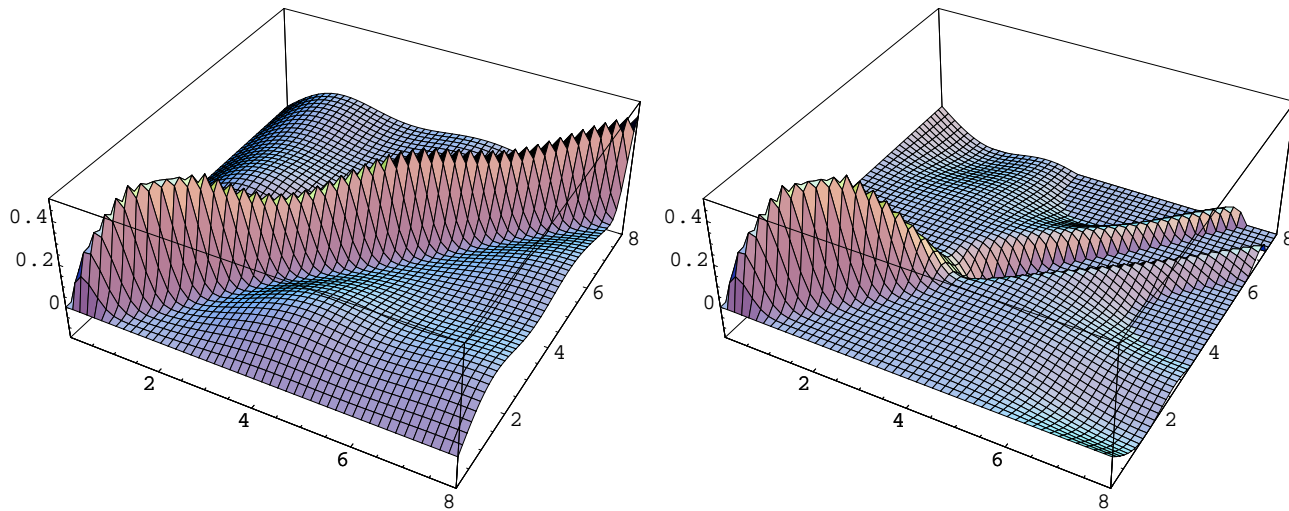
- exact result:



- expansion to 3rd order



## Expansion in 2MM



$\rho_{1,1}$  generalised Bessel vs.  $p_{1,1}$  expanded to 2nd order

- Can we find an exact result as in the 1MM?

## Reminder: Gaps in 1MM

$$E_0(s) = \frac{1}{\mathcal{Z}} \prod_{i=1}^N \int_{s^2}^{\infty} dx_i x_i^{\nu} \prod_{f=1}^{N_1} (x_i + m_f^2) e^{-Nc_1 x_i} \Delta_N(x)^2$$

- after changing variables  $x^2 \rightarrow x$
- Vandermonde  $\prod_{i>j}(x_i - x_j)$  **invariant under shift**

$\Rightarrow$   $E_0(s) \sim e^{-s^2} \times \mathbf{Z}$  **Gap = new  $\mathbf{Z}$  with**

- shifted masses  $m_f^2 \rightarrow m_f^2 + s^2$ ,
- $\nu$  extra masses  $s^2$
- “effective  $\nu = 0$ ” [Damgaard, Nishigaki, Wettig 98]

**Result**  $E_0(s) \sim e^{-s^2} \det_{f,k}[L_{N+k}(m_f^2 + s^2)] / \Delta(\{m^2 + s^2\})$   
( $\nu = 0$ )

- same strategy for  $p_k$ : shift jpdf of  $k$ -th eigenvalue  
[Damgaard, Nishigaki 01]

## Exact Results for 2MM

- partial results for objects **totally symmetric in  $\{y\}$**   
→ can replace  $\det [I_\nu(x_i y_j)]$  by diagonal  $\prod_i I_\nu(x_i y_i)$

⇒ **Gaps of  $x$  only,  $k$ -th individual  $x$ -eigenvalues nontrivial** due to presence of  $N_2$   $\mathcal{D}$ 's of  $y$ 's

- **example 1st gap:**

$$E_{0,0}(s, t = 0) \sim \prod_{i=1}^N \int_s^\infty dx_i \int_0^\infty dy_i \mathcal{P}_{jpdf}$$

- Shift easy as for 1MM?

## How to do 1st Gap in 2MM

$$E_{0,0}(s, t = 0) \sim \int_s^\infty dx \exp \text{mass1 } \Delta(x) \int_0^\infty dy \exp \text{mass2 } \Delta(y) \det I_\nu(xy)$$


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1.)  $\int_0^\infty dy$  symmetric: **replace**  $\det I_\nu(xy) \rightarrow \prod I_\nu(xy)$

2.) **replace:**

$$I_\nu(xy) \rightarrow \sum_j^\infty a_j L_j(x) L_j(y) \text{ and } \text{mass2 } \Delta(y) \rightarrow \det \begin{bmatrix} L_j(n) \\ L_j(y) \end{bmatrix}$$

$$\text{orthogonality: } \Rightarrow E_{0,0}(s, t = 0) \sim \prod_{i=1}^N \int_s^\infty dx \exp \text{mass1 } \Delta(x) \det \begin{bmatrix} L_j(n) \\ a_j L_j(x) \end{bmatrix}$$

3.) **det identity & shift & 2b.):**

$$E_{0,0}(s, t = 0) \sim \int_0^\infty dx \exp \det_{N+N_1} \begin{bmatrix} L_j(m') \\ L_j(x) \end{bmatrix} \det_{N+N_2} \begin{bmatrix} q_j(n) \\ L_j(x) \end{bmatrix}$$


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- compute like average of characteristic polynomials

## First Gap in 2MM

$$E_{0,0}(s, 0) \sim \frac{\det}{\Delta(m')\Delta(n)} \begin{bmatrix} K(m'_1, n_1) \dots K(m'_1, n_{N_2}) & L_{N+N_2}(m_1'^2) \dots L_{N+N_1-1}(m_1'^2) \\ \vdots & \vdots \\ K(m'_{N_1}, n_1) \dots K(m'_{N_1}, n_{N_2}) & L_{N+N_2}(m_{N_1}'^2) \dots L_{N+N_1-1}(m_{N_1}'^2) \end{bmatrix}$$

for  $N_1 \geq N_2$

- **shifted mass**  $m'^2 = m^2 + s^2$
- **Kernel**  $K(m', n) = \sum_{j=0}^{N+N_2-1} L_j(m^2 + s^2) q_j(n)$
- **generalised shifted Laguerre**  $q_j(n) = \sum_{l=0}^j \frac{1}{(1-\tau)^l} L_l^\nu(n^2) L_{j-l}^{-\nu-1}(s^2)$

$$1\text{MM} \lim_{\mu_1 \rightarrow \mu_2} \tau(\mu_1, \mu_2) = 0 : \quad q_j(n) = L_j(n^2 + s^2)$$

- for  $N_2 > N_1$  replace  $L_j \rightarrow q_j$  in right half of det

# Averages of Characteristic Polynomials

$$\left\langle \prod_{f=1}^{N_1} \det[i\mathcal{D}_1 + m_f] \prod_{g=1}^{N_2} \det[i\mathcal{D}_2 + n_g] \right\rangle \quad [\text{G.A., Vernizzi 03}]$$

$$= \frac{1}{\mathcal{Z}} \prod_{i=1}^N \int dx_i dy_i w(x_i, y_i) \prod_{f=1}^{N_1} (x_i^2 + m_f^2) \prod_{g=1}^{N_2} (y_i^2 + n_g^2) \Delta_N(x^2) \Delta_N(y^2)$$

$$= \frac{\text{const.}}{\Delta_{N_1}(m^2) \Delta_{N_2}(n^2)} \det \begin{bmatrix} K(m_1, n_1) \dots & K(m_1, n_{N_2}) & P_{N+N_2}(m_1^2) \dots & P_{N+N_1-1}(m_1^2) \\ \vdots & \vdots & \vdots & \vdots \\ K(m_{N_1}, n_1) \dots & K(m_{N_1}, n_{N_2}) & P_{N+N_2}(m_{N_1}^2) \dots & P_{N+N_1-1}(m_{N_1}^2) \end{bmatrix}$$

for  $N_1 \geq N_2$

## The $k$ -th $x$ -eigenvalue in 2MM

- definition

$$p_k(x) = \int_0^x dx_1 \int_{x_1}^x dx_2 \dots \int_{x_{k-2}}^x dx_{k-1} \Omega_{jpdf}(x_1, \dots, x_k)$$

- partial jpdf of the  $k$ -th ev

$$\Omega_{jpdf}(x_1, \dots, x_k) \sim \prod_{i=k+1}^N \int_{x_k}^{\infty} dx_i \prod_{j=1}^N \int_0^{\infty} dy_j \mathcal{P}_{jpdf}$$

- **Result:**

$$\Omega_{jpdf} \sim \prod_{i=k+1}^N \int_0^{\infty} dx_i \exp \det \begin{bmatrix} L_j^1(m') \\ L_j^1(x) \end{bmatrix} \det \begin{bmatrix} q_j^1(n) \\ L_j^1(x) \end{bmatrix}$$

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## The Large- $N$ Limit: Hard Edge

- rescale  $Nx_j, Ny_j, Nm, Nn$ ,  
and  $N\delta\mu^2$  'weak non-Hermiticity'

- building blocks:

- \*  $L_N(N(m^2 + s^2)) \rightarrow I_0(\sqrt{\hat{m}^2 + \hat{s}^2})$  Laguerre

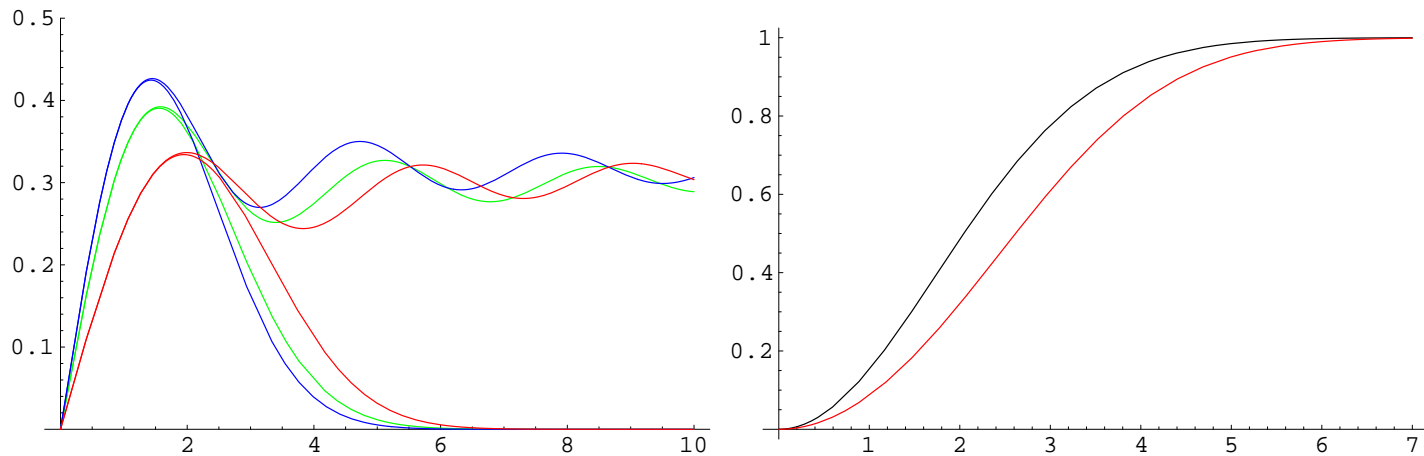
- \* 
$$q_N(n) = \sum_{l=0}^N \frac{1}{(1-\tau)^l} L_l(n^2) L_{N-l}^{-1}(s^2) \quad (\nu = 0)$$
$$\rightarrow \frac{1}{2} \int_0^1 dr e^{\frac{1}{2}r\hat{\delta}^2} I_0(\sqrt{r} \hat{n}) \frac{1}{\sqrt{1-r}} \hat{s} I_1(\sqrt{(1-r)} \hat{s}) + e^{\frac{1}{2}\hat{\delta}^2} I_0(\hat{n})$$

- tricky:  $\sum \rightarrow \int : L_0^{-1} = 1$

(check:  $\lim \hat{\delta} \rightarrow 0: \rightarrow I_0(\sqrt{\hat{n}^2 + \hat{s}^2})$  Sonine identity)

$\Rightarrow$  Kernel asymptotic =  $\sum_j L_j q_j$  & Gap

# Examples



- **left:**  $\rho_1(x)$  and  $p_1(x) \sim \partial q(\hat{n})$  for  $N_1 = 0$ ,  $N_2 = 1$   
( $\hat{\delta} \rightarrow 0$ : 1MM with 1 mass,  $\hat{\delta} \rightarrow \infty$ : 1MM quenched)
- **right:** gap  $1 - E_{0,0}(s, 0)$  comparing  
 $N_1 = 0$ ,  $N_2 = 2$  and  $N_1 = 1 = N_2$  at fixed mass and  $\hat{\delta}$

## Remarks

- for  $N_2 = 0$  integration over all  $\{y\}$ :  
**reduces 2MM to 1MM**
- **Factorisation into 1MM densities & indiv ev.:**  
for  $N_1 = N_2 = 0$  and **parameter  $\mu_j$  NOT scaled**  
with  $N$  one of the kernels converges to the weight:

$$M(x, y) \rightarrow \sum_{j=0}^{N=\infty} \chi_j(x) \hat{\chi}_j(y) = w(x, y)$$

$$\sum_{j=0}^{\infty} (1 - \tau)^j L_j(x^2) L_j(y^2) \sim I_0(xy)$$

= Mehler formula for **Hermite or Laguerre**

(compare: for scaled  $N\mu^2$ :  $\left[ \int_0^1 dt e^{-t^2 \alpha^2} J_\nu(tx) J_\nu(ty) \right]$ )

## Open Problems

- asymptotic large- $N$  rigorous and **universality proof**
- complete the solution of 2MM:  
**mixed individual eigenvalues** of  $x$ - and  $y$ -type
- **2MM with complex eigenvalues:**  
exact individual eigenvalues, partial quenching?